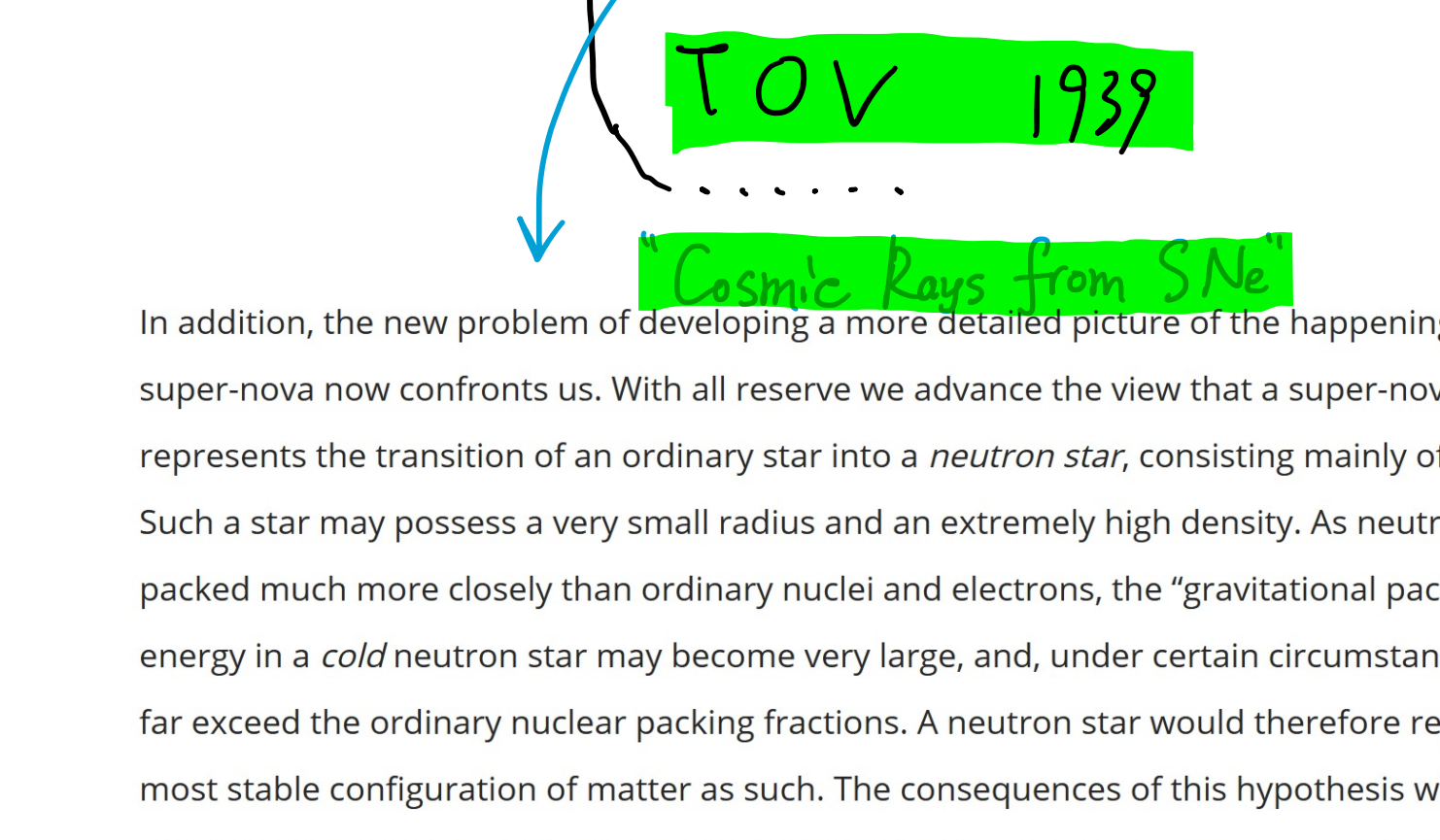


1. PSR chronicles:

是我自己加的

(1.1) Prehistoric (<1967)



In addition, the new problem of developing a more detailed picture of the happenings in a super-nova now confronts us. With all reserve we advance the view that a super-nova represents the transition of an ordinary star into a *neutron star*, consisting mainly of neutrons. Such a star may possess a very small radius and an extremely high density. As neutrons can be packed much more closely than ordinary nuclei and electrons, the "gravitational packing" energy in a cold neutron star may become very large, and, under certain circumstances, may far exceed the ordinary nuclear packing fractions. A neutron star would therefore represent the most stable configuration of matter as such. The consequences of this hypothesis will be developed in another place, where also will be mentioned some observations that tend to support the idea of stellar bodies made up mainly of neutrons.

Pacini 1967 "Energy Emission from a NS"

NS → Crab Nebula (not obs) (observed)

$$R_{NS} \sim 10^6 \text{ cm} \Rightarrow B \sim 10^8 \sim 10^9 \text{ G}$$

$$B_0 \sim 1 \sim 10 \text{ G}$$

$$R_0 \sim 10'' \text{ cm}$$

MHD waves

Relativistic particles

[没说射电强...]

No dedicated searches on radio active NS

Until Bell & Hewish (Hewish et al. 1968) 1967.7

(Pulsar = Pulse + Quasar)

Fishman et al. 1969: Obs of pulsed hard X-ray 1967.6

from NP 0532

气球

[Charles Schisler 美国中年科学家 1967 → Crab Pulsar]

(1.2) Hellas (1968-1973)

• Simple & Visual ideas.

最初 & 最简单: a magnetized ball

rotating in a vacuum.

main properties (v)

起准 TOA → rotating NS

Radio 供能

(Gunn & Ostriker 1969, 1970)

Estimation of B:

$$\vec{\Omega} \cdot \vec{\Omega} = \Omega^2$$

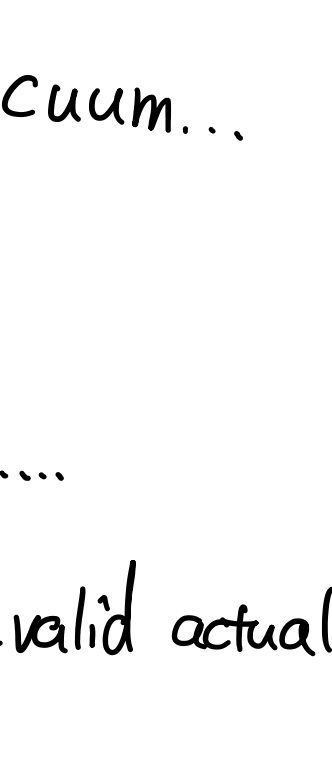
$\Omega^2 < 0$ (for most PSR)

"Spindown"

$$-I \Omega \cdot \dot{\Omega} = \frac{1}{6} \frac{B_0^2 \Omega^4 R^6}{c^3} \sin^2 \chi$$

自旋能损

磁极辐射



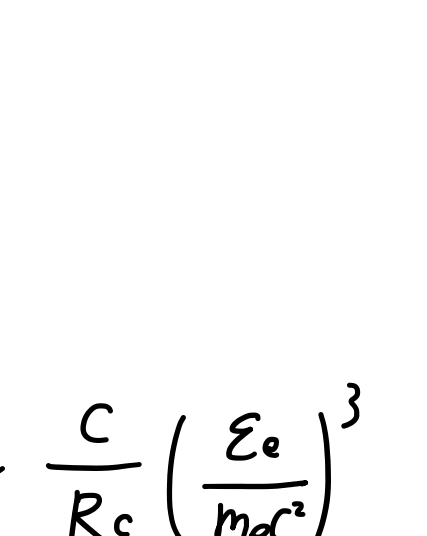
$$\Rightarrow B_{12} = (P \cdot \dot{P}_{15})^{1/2}$$

[Validity: Crab PSR: $P = 33 \text{ ms}$, $\dot{P}_{15} = 420$

$$I \Omega \cdot \dot{\Omega} \approx (2 \times 10^{33} \text{ g}) \cdot (10^6 \text{ cm})^2 \cdot (33 \text{ s}^{-1}) \cdot \left(\frac{420 \times 10^{-15}}{(33 \times 10^3)^2} \right)$$

$$\sim 10^{38} \text{ erg s}^{-1}$$

与 Crab Nebula 光学观相吻合



$$\tau_D = P / 2\dot{P} \approx 10^3 \text{ yr.}$$

(1968-1954)

Rotating Magnetic field ⇒ Electric Field

$$E \sim (v_c) \cdot B, \quad v_{\max} \sim \left(\frac{\Omega R_{NS}}{c} \right) \cdot B_0 \cdot R_{NS}$$

SI → Gauss

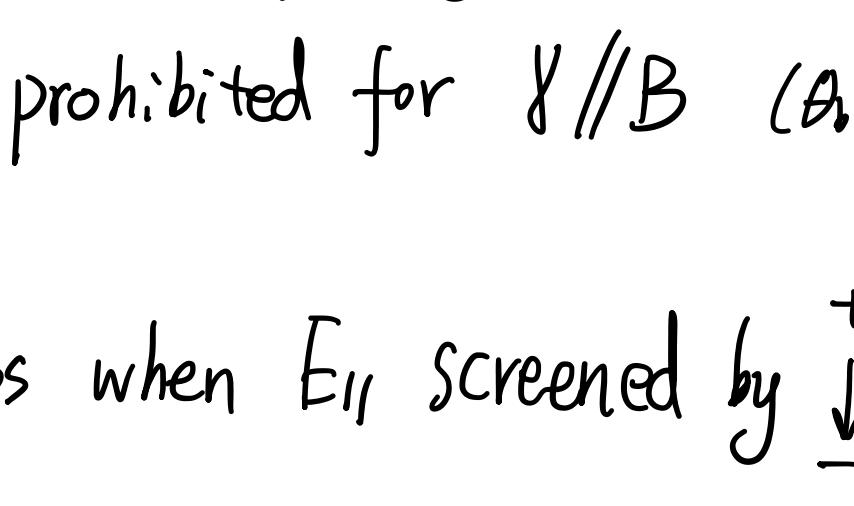
$$\approx \frac{10^{31} 10^6}{10^{10}} \cdot 10^8 \cdot 10^6 (10^2)$$

$$\sim 10^{20} \text{ V}$$

Cosmic Rays

RVM model 1969

(Radhakrishnan & Cooke)



However, PSR couldn't surrounded with vacuum...

Sturrock 1971: $\gamma + (B) \rightarrow e^+ + (B)$

magnetosphere

[$-I \Omega \dot{\Omega} = \text{Edipole}$ is invalid actually...]

how this process actually happens?

First, in vacuum rotating NS:

internal

Based on

B, G, I 1983

PSR MAGNETO

(其实不一定这样磁化)

我暂且这么想

(所以跟 magnetization 有关?)

$$\vec{J} = \sigma \left[\vec{E} + \left(\frac{\vec{\Omega} \times \vec{r}}{c} \times \vec{B} \right) \right]$$

$$0 \rightarrow \infty \Rightarrow \vec{E} + \left(\frac{\vec{\Omega} \times \vec{r}}{c} \times \vec{B} \right) = 0$$

$$\vec{E}_{||} \sim \frac{\Omega R}{c} B$$

$$\vec{E}_{||} \sim \frac{\Omega R}{c} B$$

$$\downarrow$$

$$\text{加速 } e^- \Rightarrow E_e \gg m_e c^2$$

$$e^- \text{ 沿 } \vec{B} \text{ 运动}$$

$$\text{Curvature Radiation: } \omega_{\text{cur}} \sim \frac{c}{R_c} \left(\frac{E_e}{m_e c^2} \right)^3$$

$$\text{对于电子: } \frac{dE_e}{dt} = e E_{||} - \frac{2}{3} \frac{e^2}{R_c} \left(\frac{E_e}{m_e c^2} \right)^4$$

$$\Rightarrow E_{\text{max}} \sim \left(\frac{R_c E_0}{e} \right)^{1/3} m_e c^2 \sim (10^7 - 10^8) \text{ MeV}$$

photon θ_b

$$w = \frac{3\sqrt{3}}{16\sqrt{2}} \frac{e^3 B \sin \theta_b}{\hbar m_e c^3} \exp \left(-\frac{8}{3} \frac{B_{\text{cr}}}{B \sin \theta_b} \frac{m_e c^2}{E_{\text{ph}}} \right)$$

(γ -B 反应发生概率) $\theta_b \uparrow, w \uparrow$

$$\text{其中 } B_{\text{cr}} = \frac{m_e^2 c^3}{e \hbar} \approx 4.4 \times 10^{13} \text{ G, } \hbar \omega_{\text{cur}} = m_e c^2$$

mean free path estimation:

$$\sin \theta_b^* \approx \frac{8}{3} \frac{B_{\text{cr}} m_e c^2}{B E_{\text{ph}}}$$

$$l_r = \left(\frac{R_c}{\Lambda} \right) \sin \theta_b^* \quad (\text{根据 field line 和 NS 半径})$$

$$\Lambda \sim 20, l_r \ll R_c$$

* γ -B 过程 prohibited for $\gamma \parallel B$ ($\theta_b = 0$)

e^+ 加速 γ

stops when $E_{||}$ screened by $++++$

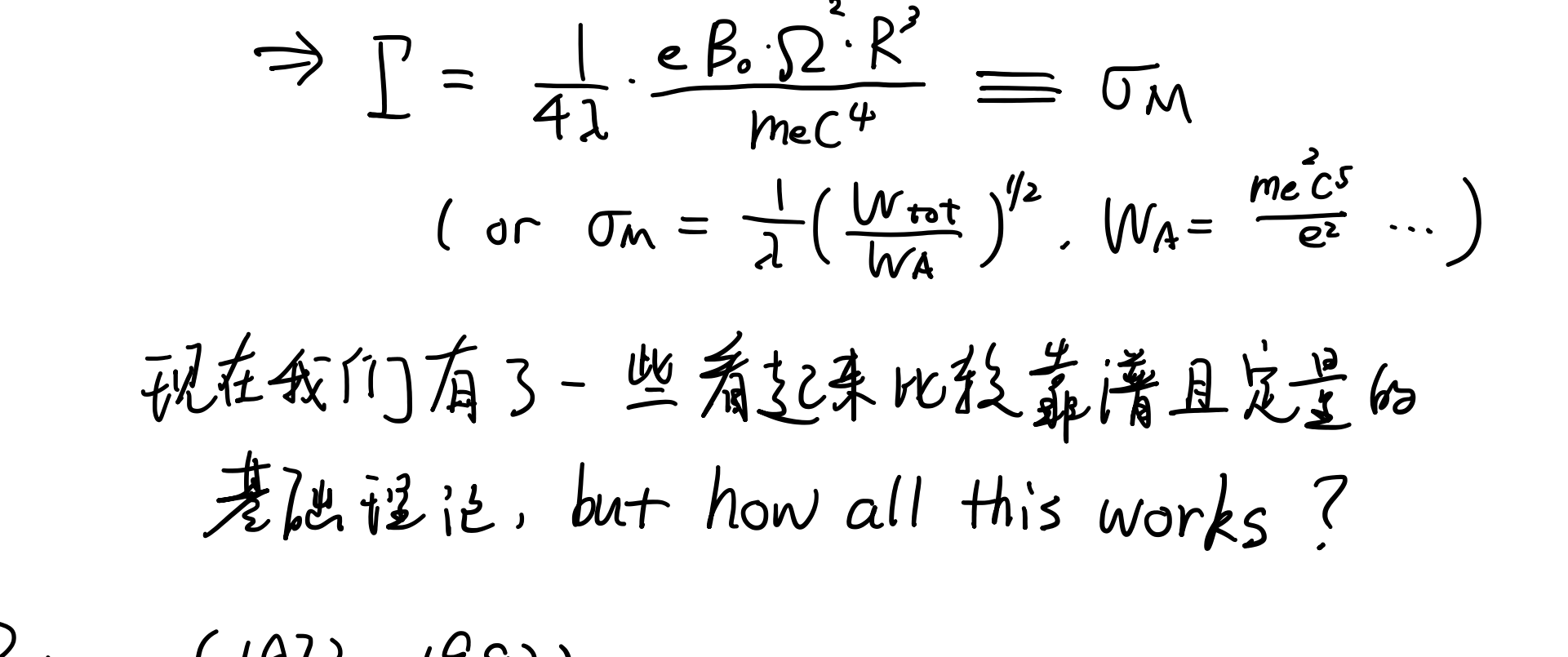
(* PSRs should be positron sources)

e^+ plasma rigidly rotates with PSR (like Earth & Jupiter's magnetospheres)

but of course fails in large distances

• when co-rotating velocity = c, $\Omega \cdot R_L = c$

$$\Rightarrow R_L = c / \Omega$$



To explain observation: kinds of radiation cone models ...

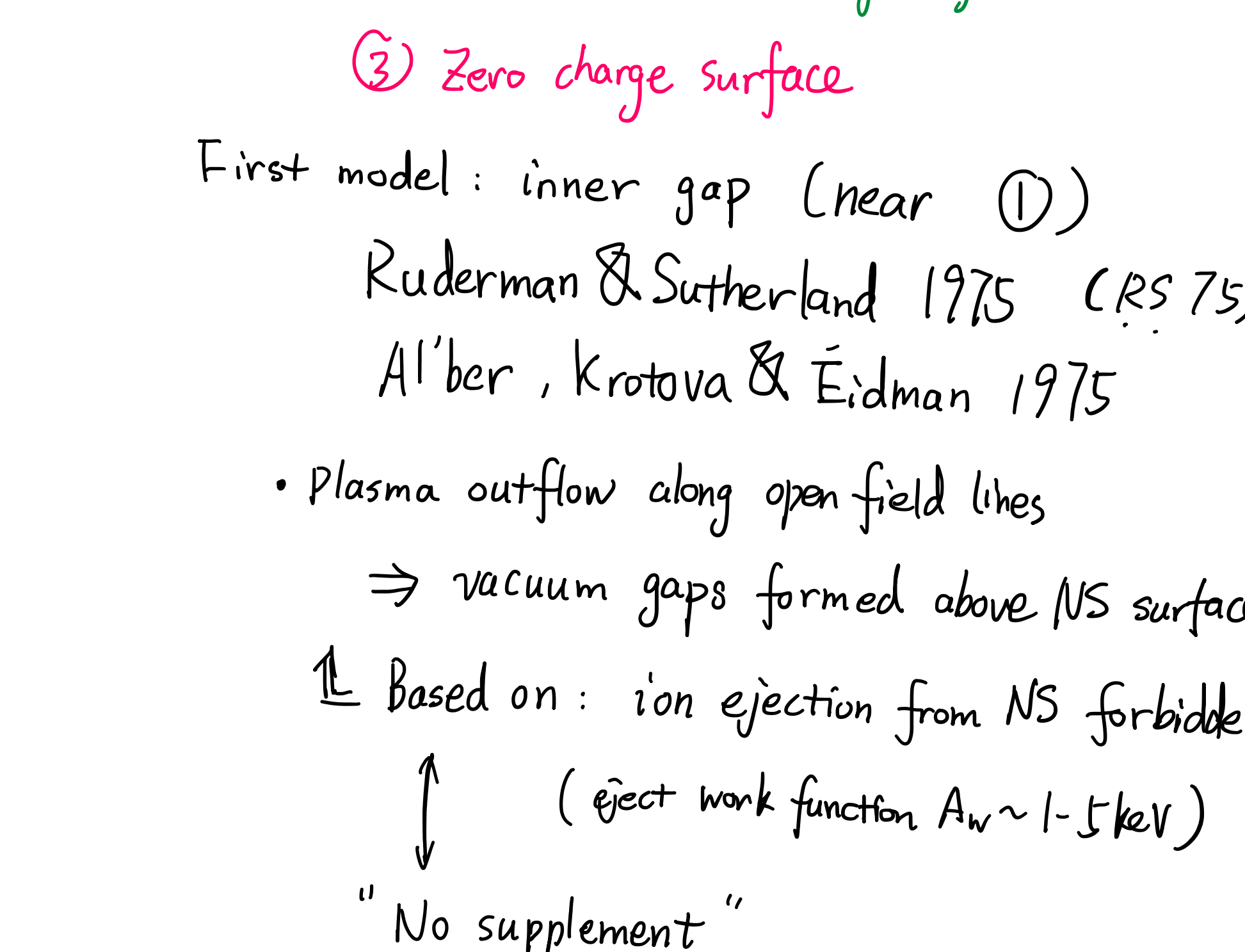
A classical one: "Hollow cone Model" (RVM 1969)

Oster & Scharlemann (1976)

Why Cone? ⇒ curvature radiation

along open field lines

Why hollow?



Measure/Define Magnetosphere parameters:

3 most important

electric charge density ρ_{AJ}

multiplication factor λ

magnetization para. σ_m

(1) ρ_{AJ} : Goldreich & Julian 1969

$$\nabla \cdot \vec{E} + \vec{\beta} \times \vec{B} = 0$$

$$\vec{\beta} = k \vec{B} + (\Omega r / c) \sin \theta \cdot \hat{\phi}$$

$$\rho_{AJ} = \frac{\nabla \cdot \vec{E}}{4\pi} = -\frac{\Omega \cdot \vec{B}}{2\pi c} \left[1 - \left(\frac{\Omega \times \vec{r}}{c} \right)^2 \right]^{-1}$$

if

$$\approx -\frac{\Omega \cdot \vec{B}}{2\pi c} \quad (\text{inner magnetosphere})$$

current

$$A \text{ balanced magnetosphere. } n_{AJ} = \frac{|\rho_{AJ}|}{e} \sim 10^{12} \text{ cm}^{-3}$$

exist

$$\text{typical current density: } j_{AJ} = \rho_{AJ} \cdot c$$

(2) $\lambda = \frac{n_e}{n_{AJ}}$

(secondary particle exceeds ρ_{AJ})

(3) $\sigma_m = \frac{1}{4\lambda} \frac{e B_0 \Omega^2 R^3}{m_e c^4}$ (Michel 1969)

$$\text{recall } W_{\text{tot}} = -I \Omega \dot{\Omega} = \frac{1}{6} \frac{B_0^2 \Omega^4 R^6}{c^3} \sin^2 \chi$$

when W_{tot} totally turn into particles' out flowing energy.

$$W_{\text{tot}} = N \cdot (I m_e c^2)$$

$$= (\lambda n_{AJ} \cdot \pi R_0^2 \cdot c) \cdot (I m_e c^2)$$

$$\Leftrightarrow \frac{1}{6} \frac{B_0^2 \Omega^4 R^6}{c^3} (\sin \theta)^2 = \lambda \cdot \frac{e B_0 \Omega^2 R^3}{2\pi c e} \cdot \pi \frac{\Omega R}{c} \cdot R^2 \cdot c \cdot I m_e c^2$$

$$\Rightarrow I = \frac{1}{4\lambda} \frac{e B_0 \Omega^2 R^3}{m_e c^4} \equiv \sigma_m$$

$$(\text{or } \sigma_m = \frac{1}{\lambda} \left(\frac{W_{\text{tot}}}{W_A} \right)^{1/2}, W_A = \frac{m_e c^2}{e} \dots)$$

现在我们有了一些看起来比较靠谱且定量的

理论基础, but how all this works?

(1.3) Rome (1973-1983)

first rigorous laws on

secondary plasma generation

magnetosphere structure

wind

.....

(a) secondary plasma generation: we need $E_{||}$

$$\text{平衡磁层: } P = \rho_{AJ}, E_{||} = 0$$

对于 PSR 磁层, 三类主要的边界 ⇒ 可能容易

① NS surface ② Light Cylinder ③ Zero charge surface

First model: inner gap (near ①)

Ruderman & Sutherland 1975 (RS 75)

Alber, Krotova & Eidman 1975

• Plasma outflow along open field lines

⇒ vacuum gaps formed above NS surface

Based on: ion ejection from NS forbidden

(eject work function $A_w \sim 1-5 \text{ keV}$)

"No supplement"

$E_{||}$ ⇒ potential drop in a certain height

解 Poisson Eq (参见 RS 1975 附录 III)

可知 Polar Cap 能产生的 maximum 电压:

$$\Delta V_{\text{max}} = \frac{B_0 R \Omega^2}{2 c^2}$$

$$\text{secondary plasma 产生: } H_{\text{rs}} \approx 10^4 B_{12}^{-1/2} P^{3/4} R_7^{3/4} \text{ (cm)} \quad (R_7 = \frac{R_0}{10 \text{ cm}} \text{ polar cap})$$

$$\psi_{\text{rs}} \approx 6 \times 10^{12} B_{12}^{-1/2} P^{1/4} R_7^{1/4} \text{ (v)}$$

"turn off": $\psi_{\text{rs}} = \Delta V_{\text{max}}$

(或 $H_{\text{rs}} = R_0$)

One question of RS: ion ejection forbidden?

Later estimation: $A_w \sim 100 \text{ eV} \ll 1 \sim 5 \text{ keV}$ for NS

替代方案: Fowler & Arons & Scharlemann 1977

Arons & Scharlemann 1979

.....

free-flow

来不及推完了.....

PSR Eq.

Current Problem

Wind Problem

PIC 时代 ...